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Magnetic Fields in the Universe V Cargèse, Corsica

INTRODUCTION

Synchrotron radiation carries the imprint of the magnetic field at the point of origin and along the propagation path.

Linear polarisation vector Amplitude of polarisation Angle of polarisation

P = Q + iU $|P| = \sqrt{Q^2 + U^2}$ $\theta = (1/2) \tan^{-1}(U/Q)$

Faraday rotation Faraday depth

$$\begin{aligned} \theta &= \theta_0 + \phi \lambda^2 \\ \phi &= 0.81 \int_{\text{source}}^{\text{observer}} n_e \mathbf{B} \cdot d\mathbf{l} \end{aligned}$$

THE GRADIENT OF LINEAR POLARISATION

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LETTER

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Low-Mach-number turbulence in interstellar gas revealed by radio polarization gradients

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$$|\nabla \mathbf{P}| = \sqrt{\left(\frac{\partial \mathbf{Q}}{\partial \mathbf{x}}\right)^2 + \left(\frac{\partial \mathbf{U}}{\partial \mathbf{x}}\right)^2 + \left(\frac{\partial \mathbf{Q}}{\partial \mathbf{y}}\right)^2 + \left(\frac{\partial \mathbf{U}}{\partial \mathbf{y}}\right)^2}$$

 $|\nabla P|$ measures the rate at which the polarisation vector traces out a trajectory in the Q–U plane as a function of position on the sky.

THE GRADIENT OF LINEAR POLARISATION



Figure

- Smoothed and unsmoothed signal
- Derivative of the unsmoothed signal
 - May enhance noise in the data (Burkhart, Lazarian & Gaensler 2012)
- Derivative of the smoothed signal



Why should we smooth maps before calculating the gradient?

Robitaille & Scaife 2015 MNRAS, 451 (1), 4891



Figure: The spatial gradient of linearly polarized emission, $|\nabla P|$, at the original resolution (left) and for the smoothed Stokes Q and U maps (right). Smoothed maps are produced with a convoluted Gaussian filter having a standard deviation of 2^2 pixels. (Canadian Galactic Plane Survey (CGPS) polarised emission at 1420 MHz (Landecker et al. 2010 – Dr. Roland Kothes's poster)





Statute Language (

10



- The direct convolution of the Derivative of a Gaussian (DoG)
- The function satisfies the properties of a wavelet transform

Wavelet Transform

$$\tilde{f}(l, \mathbf{x}) = \begin{cases} \tilde{f}_1 = \frac{1}{l^2} \int \psi_1 \left[\frac{(\mathbf{x}' - \mathbf{x})}{l}\right] f(\mathbf{x}') d^2 \mathbf{x}' \\ \\ \tilde{f}_2 = \frac{1}{l^2} \int \psi_2 \left[\frac{(\mathbf{x}' - \mathbf{x})}{l}\right] f(\mathbf{x}') d^2 \mathbf{x}', \end{cases}$$

where

$$\psi_1(\mathbf{x}, \mathbf{y}) = \frac{\partial^m \phi(\mathbf{x}, \mathbf{y})}{\partial \mathbf{x}^m}$$
 and $\psi_2(\mathbf{x}, \mathbf{y}) = \frac{\partial^m \phi(\mathbf{x}, \mathbf{y})}{\partial \mathbf{y}^m}$

Original Gradient of Linear Polarisation:

$$|\nabla \mathsf{P}| = \sqrt{\left(\frac{\partial \mathsf{Q}}{\partial x}\right)^2 + \left(\frac{\partial \mathsf{U}}{\partial x}\right)^2 + \left(\frac{\partial \mathsf{Q}}{\partial y}\right)^2 + \left(\frac{\partial \mathsf{U}}{\partial y}\right)^2}$$

Multiscaled Gradient of Linear Polarisation:

$$|\nabla \tilde{\mathsf{P}}(\mathfrak{l},\mathbf{x})| = \sqrt{|\tilde{\mathsf{Q}}(\mathfrak{l},\mathbf{x})|^2 + |\tilde{\mathsf{U}}(\mathfrak{l},\mathbf{x})|^2},$$

$$\begin{split} |\tilde{Q}(l,x)| &= \sqrt{|\tilde{Q}_1(l,x)|^2 + |\tilde{Q}_2(l,x)|^2}, \\ |\tilde{U}(l,x)| &= \sqrt{|\tilde{U}_1(l,x)|^2 + |\tilde{U}_2(l,x)|^2}. \end{split}$$



Figure: From left to right: the $|\nabla \tilde{P}(l, \mathbf{x})|$ values at four different scales l= 9.6, 45.7, 153.6 and 434.4 arcmin. White lines represent maxima chains corresponding to the scale (WTMM, Arnéodo et al. 2000, European J. Phys. B, **15**, 567).



Figure: The superposition of maxima chains from scale l =22.8 to 258.3 arcmin over the map of $|\nabla \tilde{P}|$ at l = 22.8 arcmin. They represent a subset maxima chains for which the maximum value along the chain is part of outliers separated with the scale-wise CVE algorithm.

POWER SPECTRUM

- The power spectrum of an image can be calculated from its wavelet coefficients.
- Δ-variance analysis (Stutzki et al. 1998, Bensch et al. 2001, Ossenkopf et al. 2008)
- Some directional wavelets can reproduce the classical Fourier power spectrum (Kirby 2005, Robitaille, Joncas & Miville-Deschênes 2014).

$$\int |\mathbf{f}(\mathbf{x})|^2 d^2 \mathbf{x} = C_{\psi}^{-1} \int \int \frac{|\tilde{\mathbf{f}}(\mathbf{l}, \mathbf{x})|^2}{\mathbf{l}^2} d\mathbf{l} d^2 \mathbf{x}$$
(1)

$$E(l) = \int \frac{|\tilde{f}(l, \mathbf{x})|^2}{l^2} d^2 \mathbf{x} \qquad (2)$$

$$S_{P}(l) = \frac{1}{N_{x}N_{y}}\sum_{\mathbf{x}}|\nabla \tilde{P}(l,\mathbf{x})|^{2} \quad (3)$$

Comparison between the wavelet power spectrum of $|\nabla \tilde{P}|$ and the Fourier power spectrum of |P|



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POWER SPECTRUM

Comparison between the wavelet power spectrum of $|\nabla \tilde{P}|$ and the Fourier power spectrum of |P|



Figure 12. Normalised distributions of wavelet coefficients of $\langle \nabla \hat{P}(l, x) \rangle$ for the CGP8 field for all coefficients (top panel). The black lines represent scales between 6.8 and 9.13 arcmin and the blue lines present scales between 108.6 and 614.4 arcmin. The lower panel shows the Gaussian part of the distribution for scales between 6.8 and 9.1.3 arcmin.



CONCLUSION

- It is possible to extend the calculation of $|\nabla \tilde{P}|$ to multiple scales using a wavelet analysis formalism.
- \cdot Fluctuations traced by $|\nabla \tilde{P}|$ exist at larger scales on data completed with lower spatial frequencies.
- \cdot We can measure the power spectrum of $|\nabla \tilde{P}|$ using the wavelet formalism.
- This analysis will be applied to different radio polarisation data (S-PASS) and the algorithm adapted for calculation on a sphere.
- Some features ("large-scale" double jumps, coherent features across scales, ...) must be compared with simulations.