



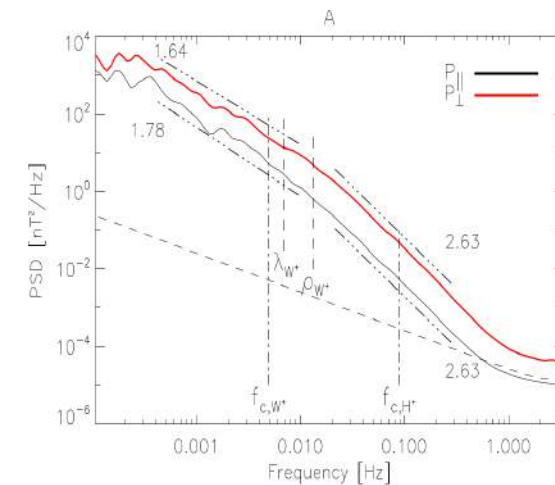
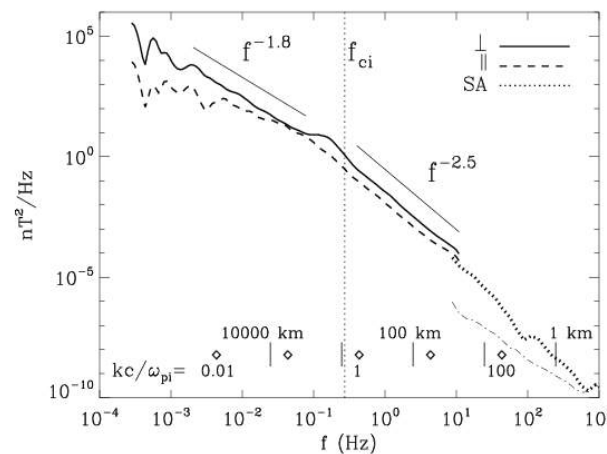
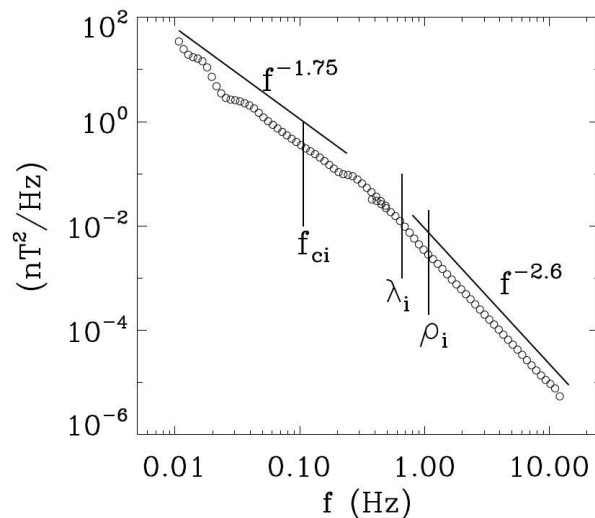
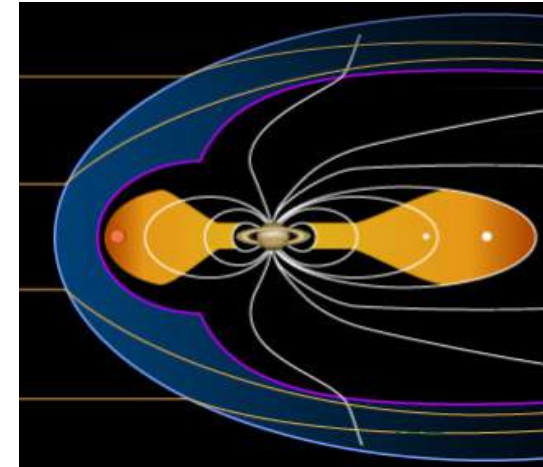
Observations of solar wind turbulence (from MHD to electron scales)

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Plasma Turbulence in the Heliosphere

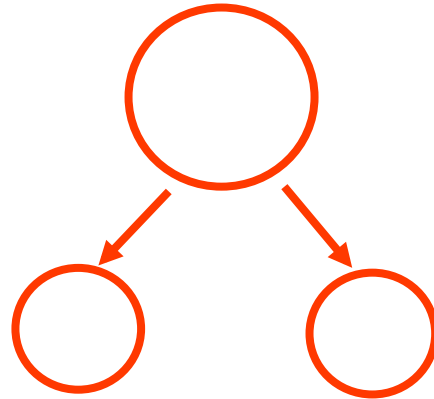
In situ measurements in the solar wind and planetary magnetospheres show omnipresence of plasma turbulence.



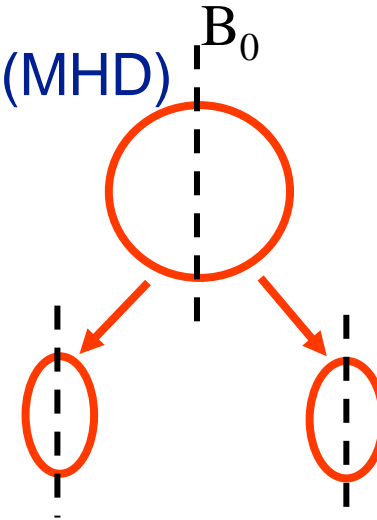
[Alexandrova et al. 2013, 2008, Von Papen et al. 2014]

Turbulence in space plasmas

hydrodynamics



plasma (MHD)



1. Presence of a mean magnetic field B_0 leads to an anisotropy of turbulent fluctuations.
2. Plasma waves: Alfvén, magnetosonic, mirror, whistlers, kinetic Alfvén waves (KAW), etc... (wave turbulence).
3. No collisions : m.f.p. ~ 1 AU.
4. In plasmas there is a number of characteristic space and temporal scales.

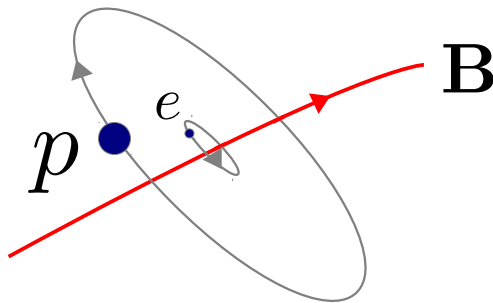
$$f_{ci}, c/\omega_{pi}, R_{Li}$$

$$f_{ce}, c/\omega_{pe}, R_{Le}$$

$$\lambda_D$$

Typical plasma scales in the solar wind

- Larmor radius ($\rho_i \sim 100 \text{ km}$, $\rho_e \sim 1 \text{ km}$) and cyclotron frequency ($\Omega_{ci}/2\pi \sim 0.1 \text{ Hz}$, $\Omega_{ce}/2\pi \sim 200 \text{ Hz}$) of a charged particle (electron or ion=proton) in a magnetic field B:



$$\rho_{i,e} = \frac{V_{\perp i,e}}{\Omega_{ci,e}} ; \Omega_{ci,e} = \frac{eB}{mc}$$

- Inertial length $\lambda_{i,e}$ (scale of the demagnetization of the particles, which is close to $\rho_{i,e}$ in plasma with $\beta = nkT/(B^2/2\mu_0) \sim 1$) and plasma frequency (ω_p):

$$\lambda_{i,e} = \frac{c}{\omega_{pi,e}} ; \omega_{pi,e}^2 = \frac{4\pi ne^2}{m_{i,e}}$$

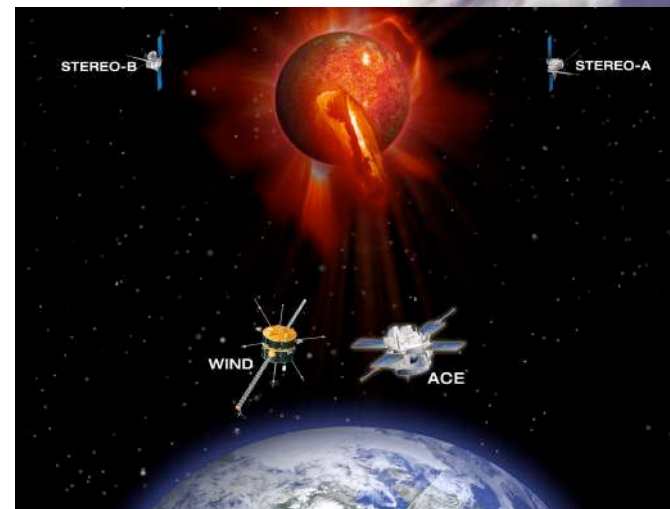
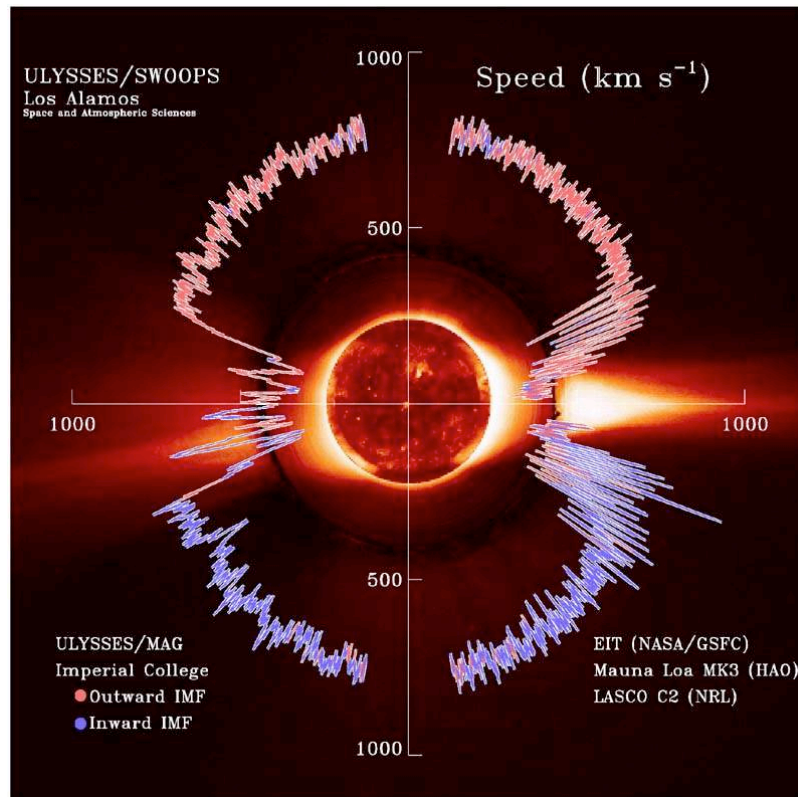
- Debye length $\lambda_D \sim 10 \text{ m}$ (sphere of influence of a given test charge in a plasma); at $L > \lambda_D$ plasma is quasi-neutral:

$$\lambda_D^2 = \frac{k_B T}{8\pi n e^2}$$

- Satellite size $\sim 1-2 \text{ m}$.

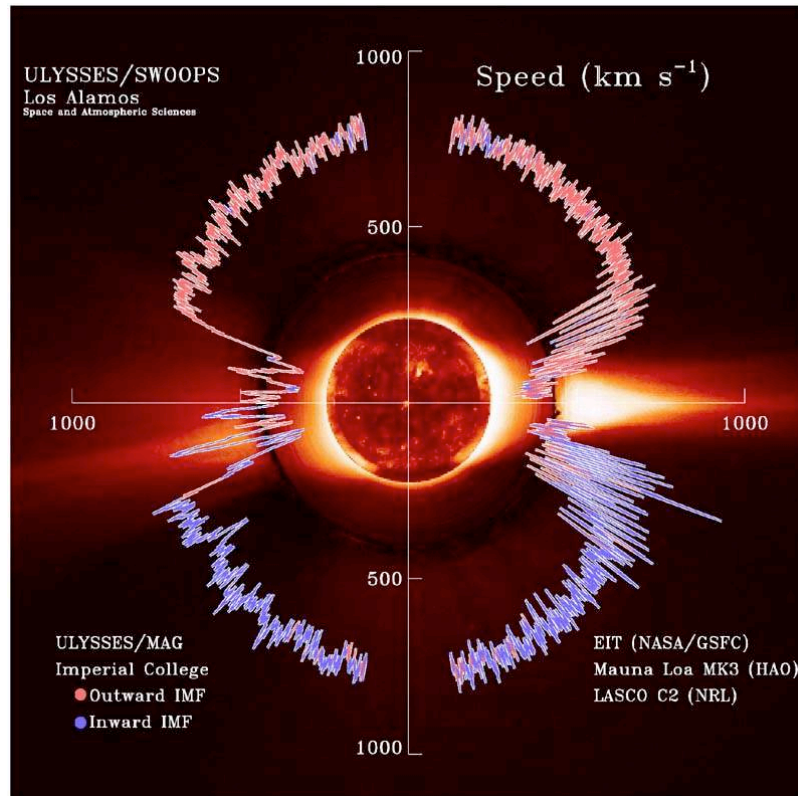
The solar wind

one of the best and the closest laboratory for astrophysical plasma turbulence.

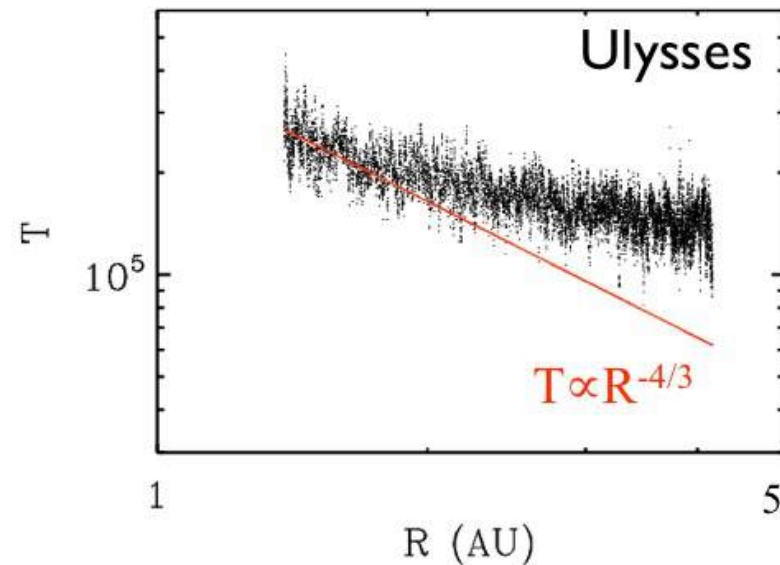


Two components, Slow and Fast streams.
Slow wind: $V = 300\text{-}400 \text{ km/s}$, $n=7 \text{ cm}^{-3}$, $T_p=2 \cdot 10^5 \text{ K}$
Fast wind: $V = 600\text{-}800 \text{ km/s}$, $n=3 \text{ cm}^{-3}$, $T_p=5 \cdot 10^5 \text{ K}$

The solar wind



[Image credit: Lorenzo Matteini]

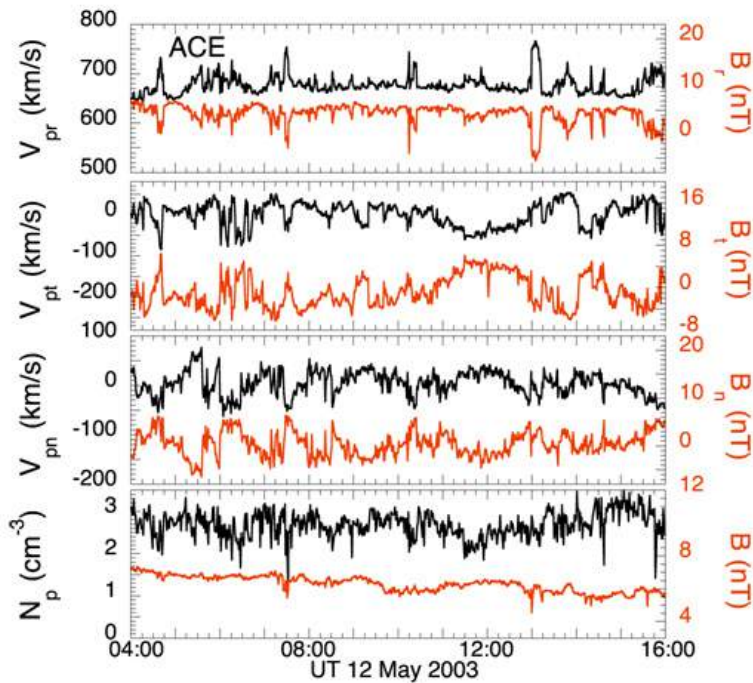


Wind temperature decays
less than adiabatic ($\sim R^{-4/3}$)

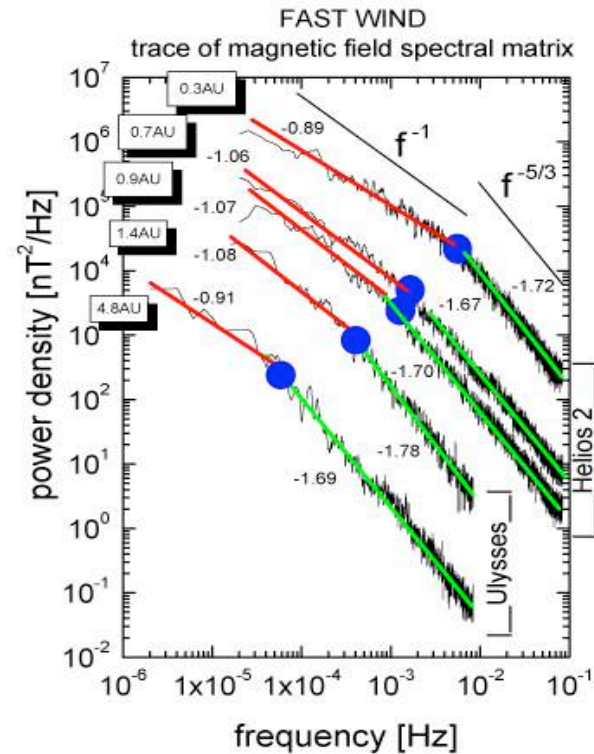
Turbulence dissipation may explain the solar wind heating [e.g. Vasquez et al. 2007; Sorriso-Valvo et al. 2007; Macbride et al. 2008; Smith et al. 2009; Cranmer et al. 2009; Marino et al. 2012; Wu et al. 2013, ...]

Solar wind Turbulence and Alfvén waves

[Gosling et al., 2009; Balcher & Davis 1977]

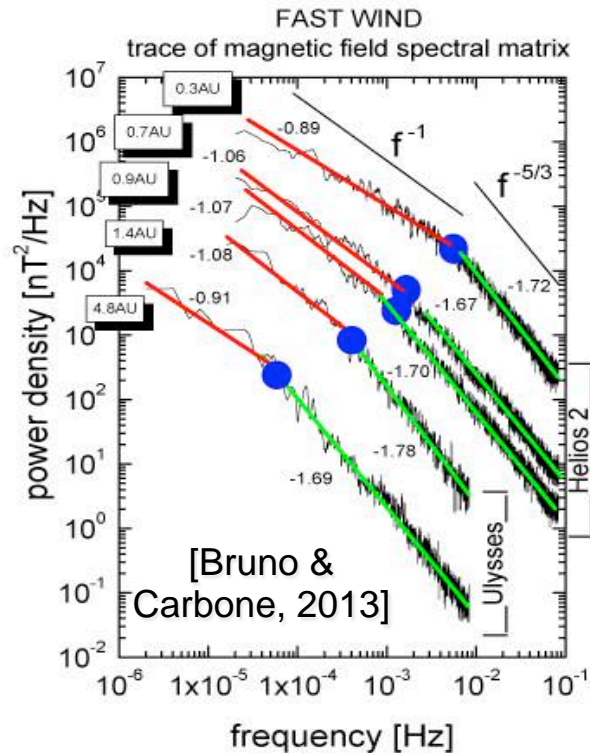


[Bruno & Carbone, 2013]



- Strong correlation between V and B fluctuations at 1 AU (Alfvén waves)
- These waves belong to f^{-1} spectral range.
- Kolmogorov turbulence at smaller scales (MHD) is observed.

Starting point of the Kolmogorov spectrum

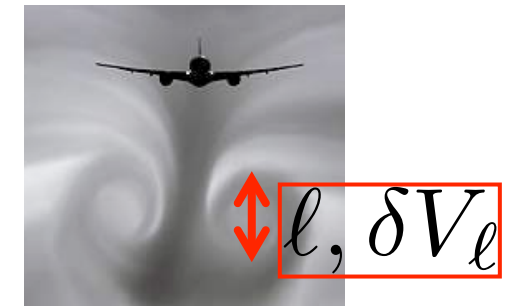
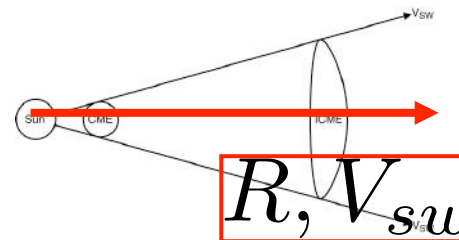


- The solar wind expansion time:

$$\tau_{exp} = R/V_{sw}$$

- The eddy-turnover time:

$$\tau_{NL} = \ell/\delta V_{\ell}$$

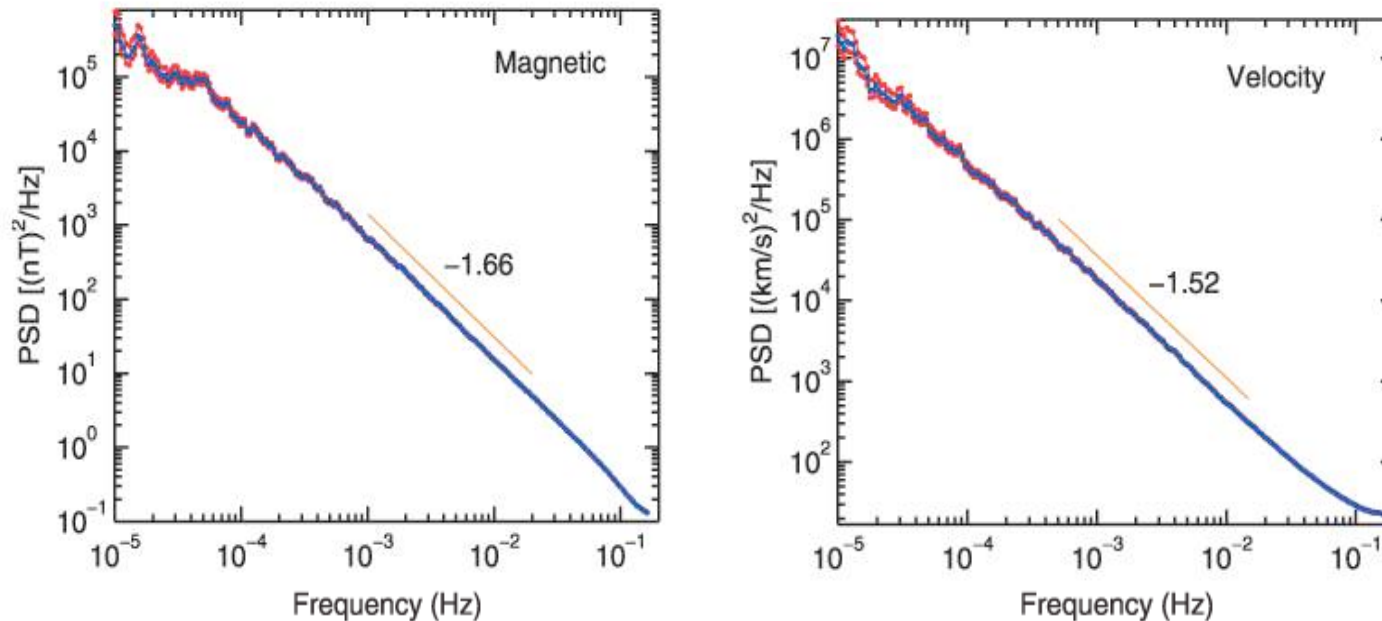


- Transition between f^{-1} and $f^{-5/3}$ spectrum corresponds to a scale where these 2 characteristic times are of the same order [Mangeney et al. 1991; Meyer-Vernet 2007]:

$$\tau_{exp} \simeq \tau_{NL}$$

Solar wind Turbulence and Alfven waves

In a case of a pure alfvenic turbulence magnetic and velocity spectra should be the same, but in the solar wind it is not the case:



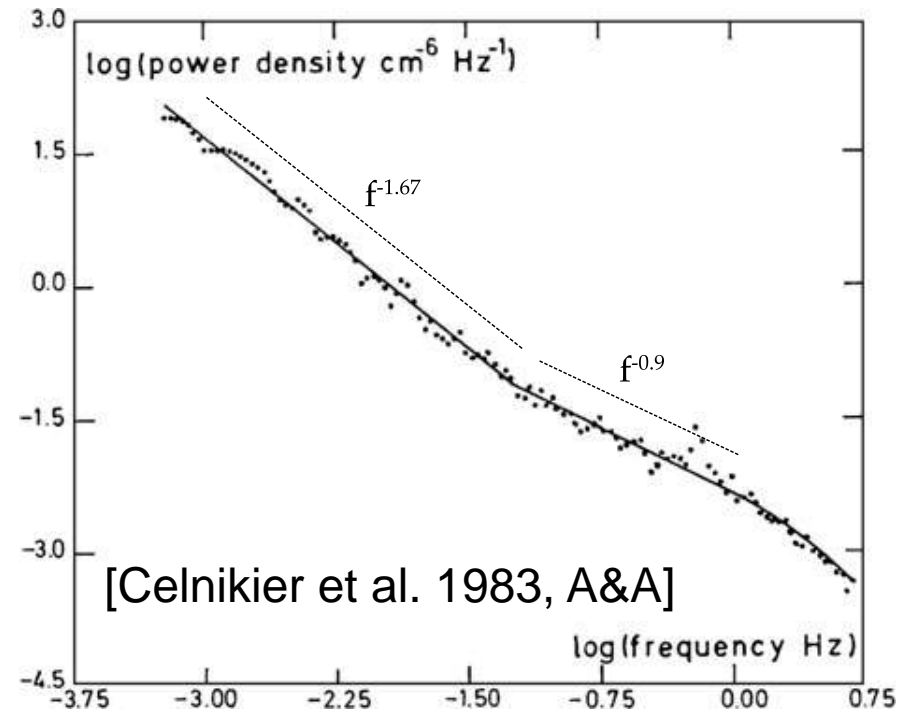
[Podesta et al., 2007; Salem 2000, PHD]

Why?

- Local dynamo process (Grappin et al., 1983) ?
- Solar wind expansion ? V-B alignment (see talk of S. Boldyrev)?
- Compressibility ?

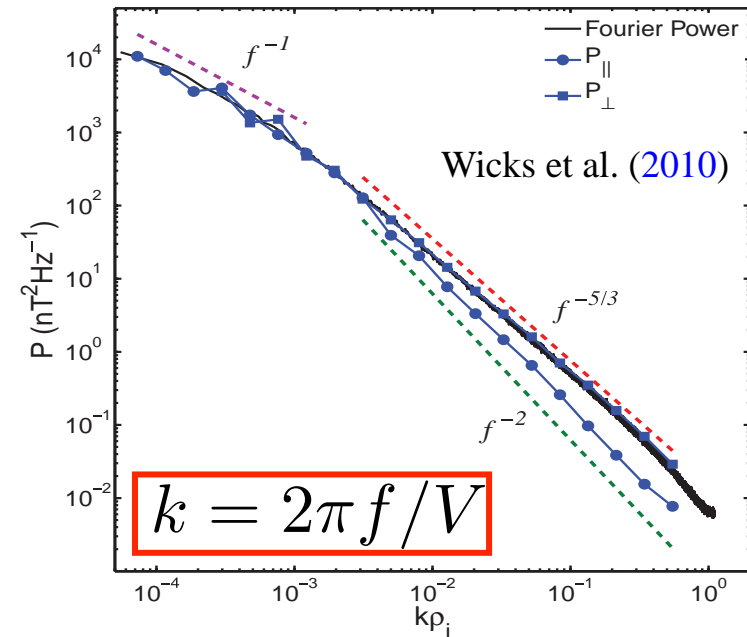
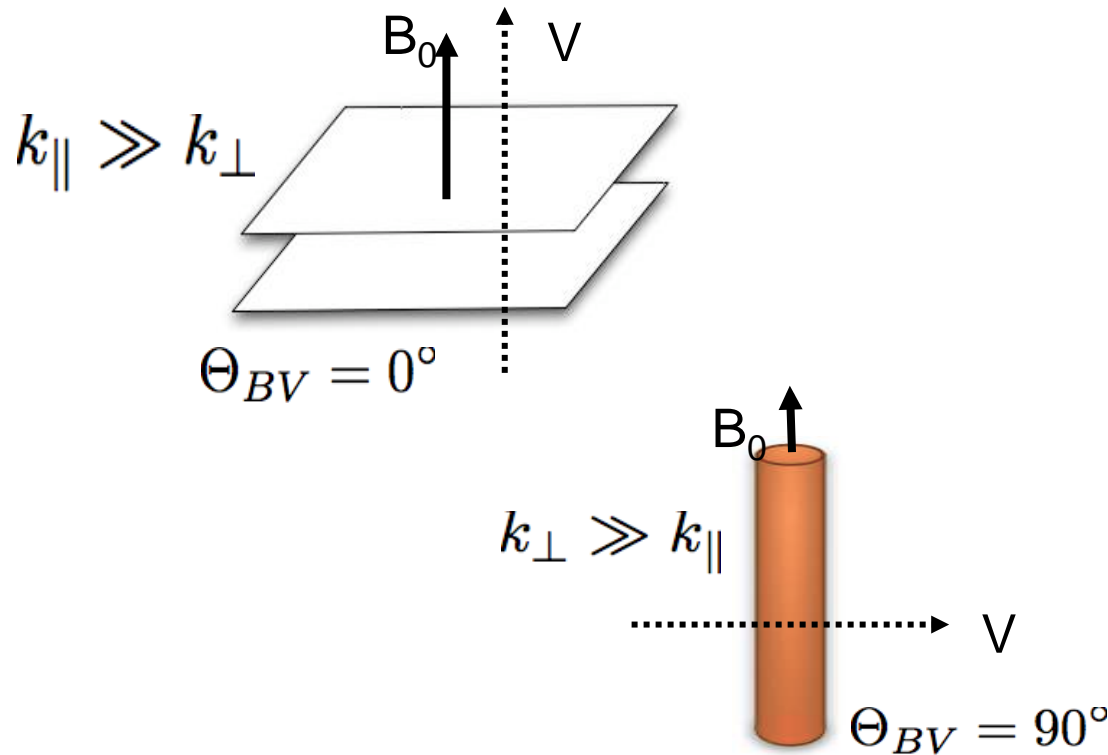
Solar wind turbulence is compressible

Spectrum of electron density fluctuations in the solar wind as measured by ISEE 1 & 2. See as well Chen et al. 2013.



Can the compressibility be the source of the non-alfvenicity of the inertial range in the solar wind turbulence?

Anisotropy of turbulent fluctuations at MHD scales



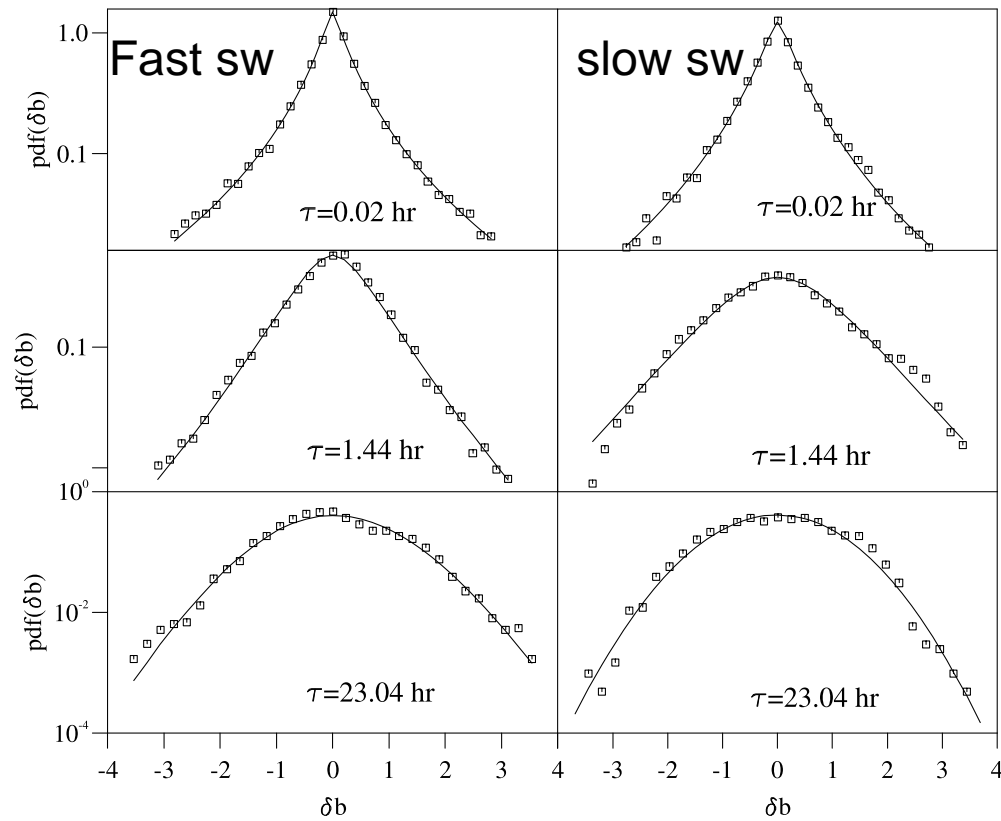
- Alfvénic turbulence of Goldreich and Sridhar, 1995 is based on the idea of a balance between linear Alfvén time (along B_0) and non-linear time (in plane perp. to B_0):

$$\tau_A = \frac{\ell_{\parallel}}{V_A} \sim \tau_{NL} = \frac{\ell_{\perp}}{\delta V_{\perp}}$$

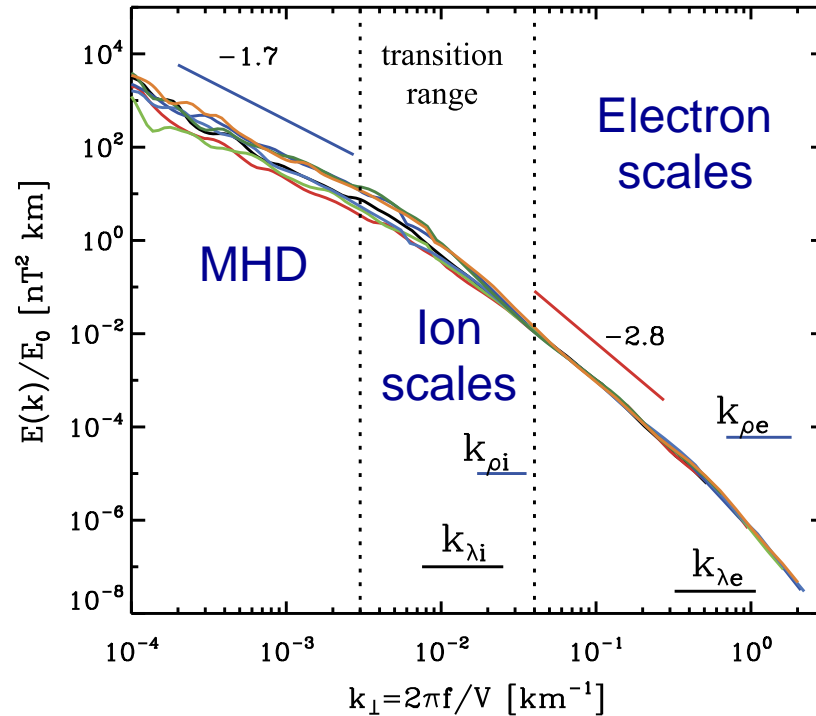
$$P(k_{\perp}) \sim k_{\perp}^{-5/3} ; P(k_{\parallel}) \sim k_{\parallel}^{-2}$$

Intermittency of turbulent fluctuations within the MHD inertial range

[Sorriso-Valvo et al. 1999, Dudok de Wit et al. 2013]



Solar wind turbulent spectrum of magnetic fluctuations at MHD-Ion-Electron scales



[Alexandrova, Chen, Sorriso-Valvo, Bale, Horbury, 2013 Space Science Rev, open access]

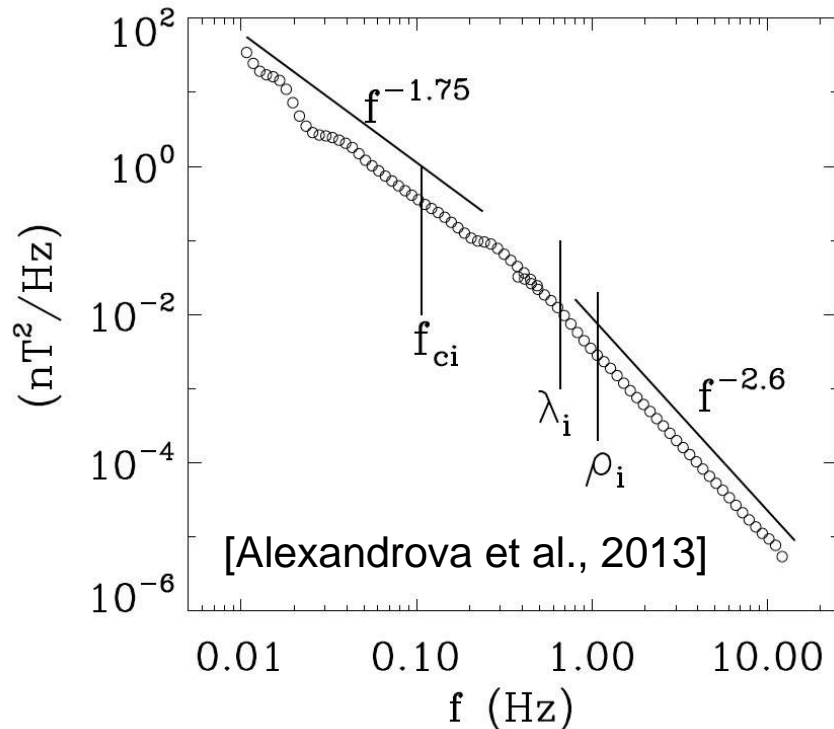
1. What is going on close to ion and electron scales?
2. Which plasma scale is responsible for the ion break?
3. Which plasma scale plays the role of the dissipation scale?
4. Physical mechanisms?
5. Nature of turbulent fluctuations : waves or strong turbulence?
6. ...

Turbulence at kinetic scales

1. Ion scales

$$f_{ci} = \frac{eB_0}{2\pi m_i c}, \quad k\rho_i \sim 1, \quad kc/\omega_{pi} \sim 1$$

Which ion scale is responsible for the break?



Time scale

$$f_{ci} = \Omega_{ci}/2\pi ; \Omega_{ci} = eB/m_i c$$

Spatial scales

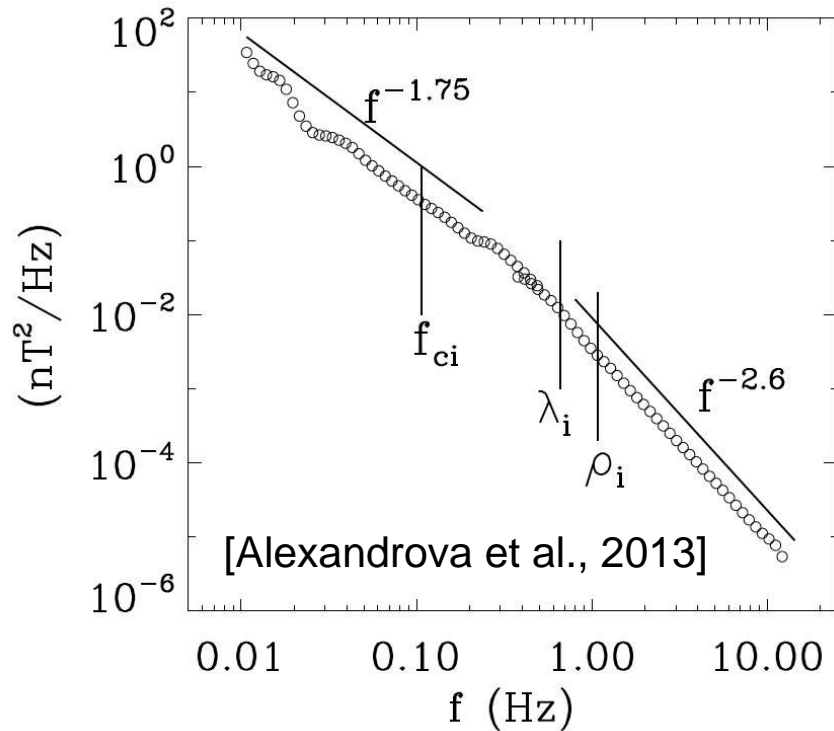
$$\rho_i = \frac{V_{\perp i}}{\Omega_{ci}} ; \lambda_i = \frac{c}{\omega_{pi}} = \frac{V_A}{\Omega_{ci}}$$

In frequency spectrum, these scales appear at Doppler shifted frequencies:

$$f_{\rho_i} \simeq \frac{V_{solar\ wind}}{\rho_i} ; f_{\lambda_i} \simeq \frac{V_{solar\ wind}}{\lambda_i}$$

- All characteristic time and spatial ion scales are observed close to the spectral break point...
- How can we distinguish between different scales?
- Important in order to understand which physical mechanisms “break the spectrum” (e.g., if it is f_{ci} => damping of Alfvén waves).

Which ion scale is responsible for the break?



- Leamon et al. 2000 : λ_i
- Schekochihin et al. 2009: ρ_i
- Perri et al. 2010 : any of the scale/combination of scales
- Bourouaine et al. 2012: λ_i
- Bruno et al. 2014: resonant k of parallel Alfvén waves
- Chen et al. 2014: beta dependent.

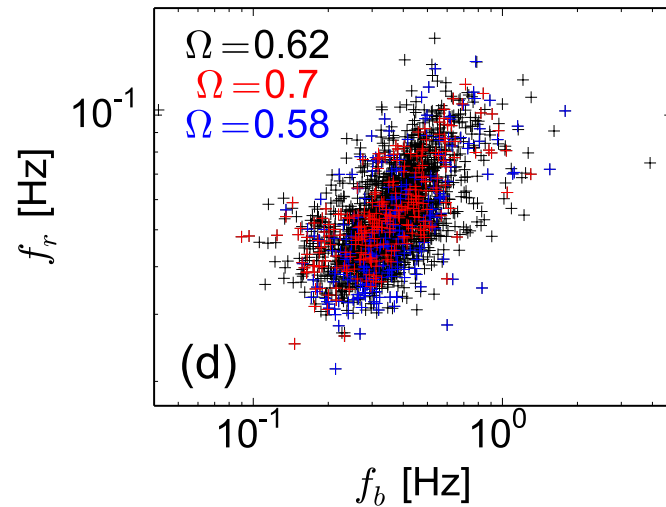
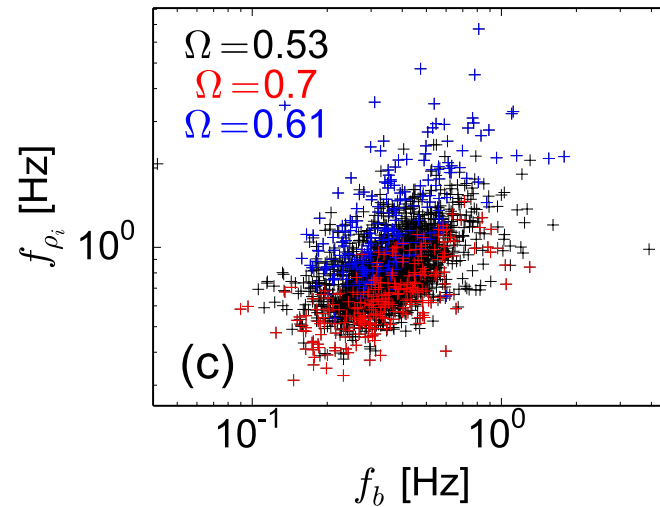
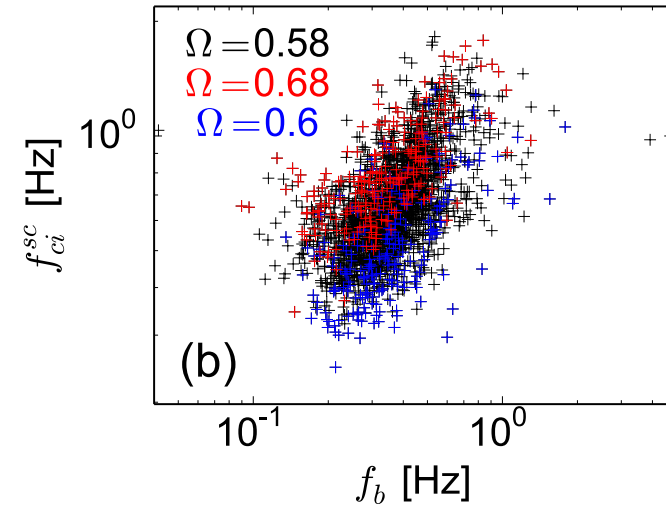
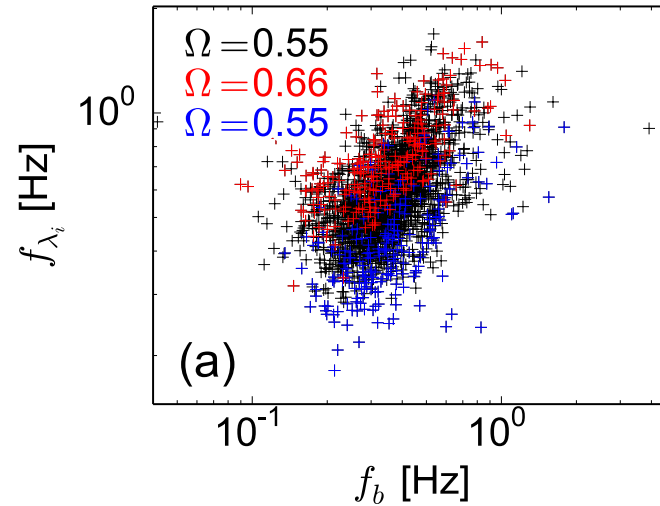
$$\beta_i = 2\mu_0 n k_B T_i / B^2 = \rho_i^2 / \lambda_i^2.$$

⇒ The largest characteristic ion scale “breaks” turbulent spectrum [Chen et al. 2014].

All scales correlates well with f_b ...

[Sonny Lion, 6 years of STEREO/MAG data]

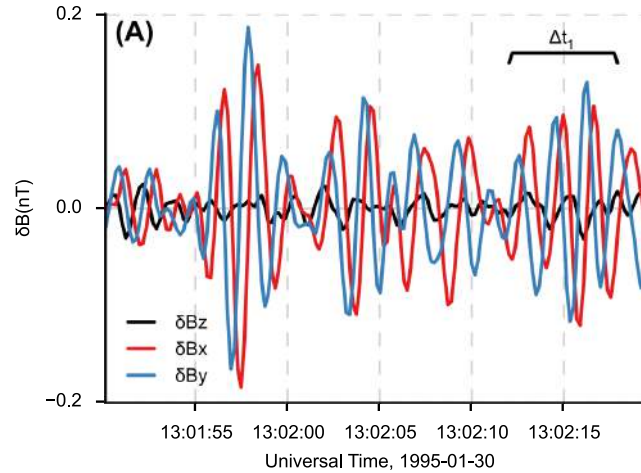
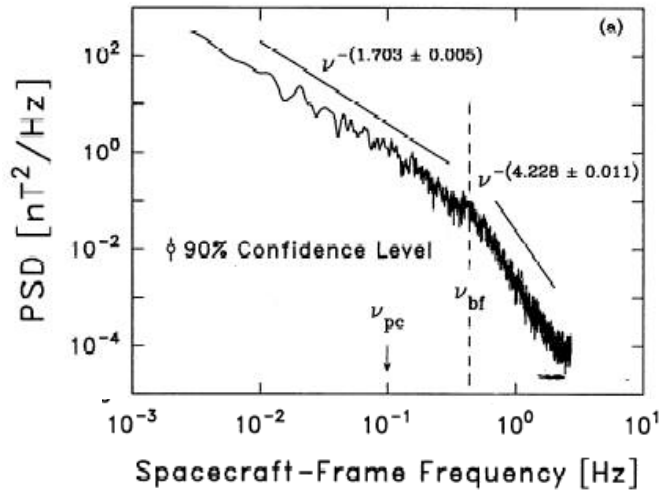
Red: high $\beta_p > 1$
Blue: low $\beta_p < 0.2$
Black: all data



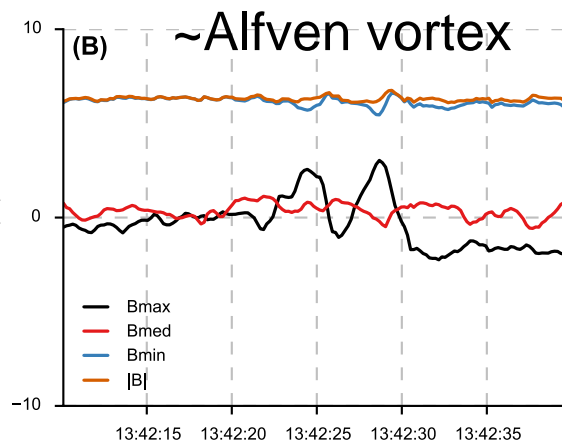
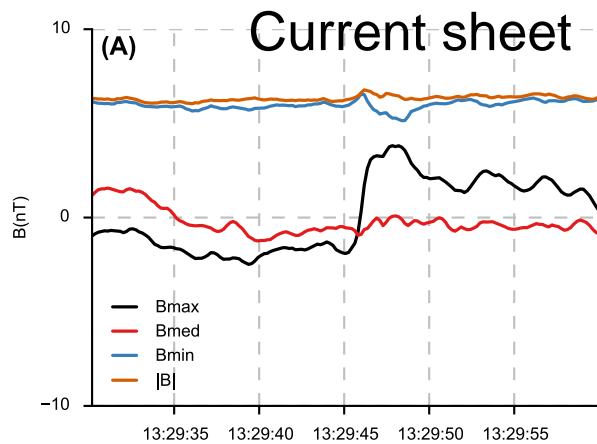
\Rightarrow No ONE scale (ONE physical phenomena) to explain the ion transition ?

Ion scales: superposition of different phenomena

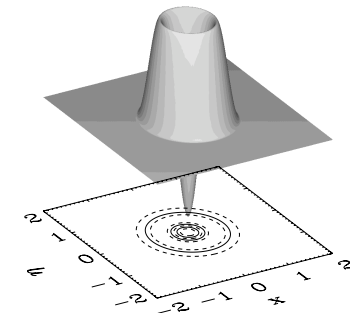
[Leamon et al. 1998]



Monochromatic Alfvén waves at $\text{freq} \sim f_{ci}$ with $k_{||}$ (generated by AIC instability).



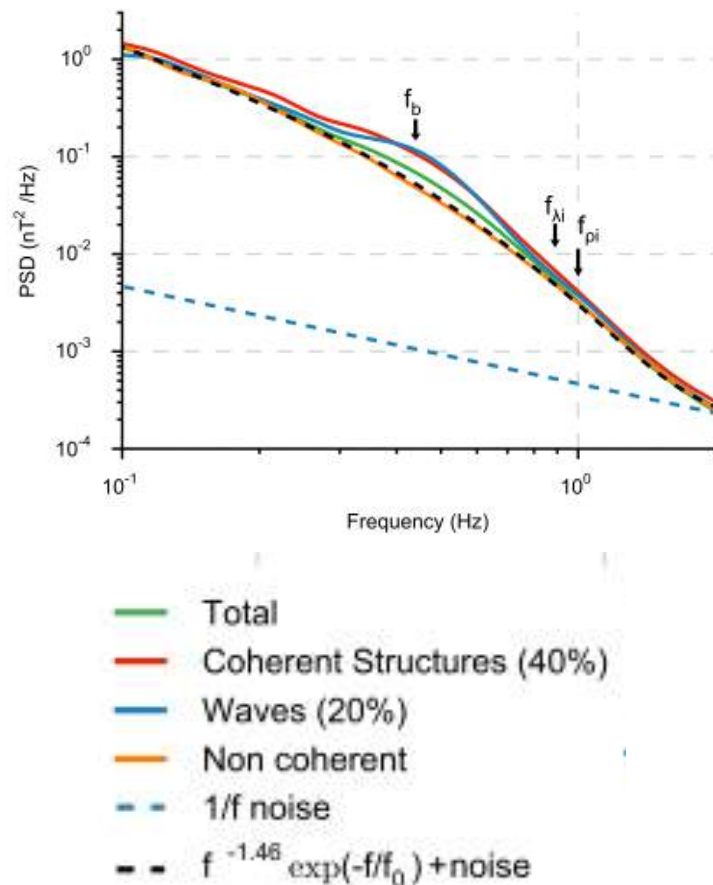
Localized spatial structures with k_{perp} at scale \sim ion Larmor Radius



[Lion et al, 2015, submitted to APJ]

Solar wind spectrum around ion scales

(conclusions based on a case of a fast sw stream)



It consists of

- Alfvén Ion Cyclotron waves (with k_{\parallel})
- Coherent structures (with k_{perp})
- Non coherent signal, which can be described by

$$E_B \sim f^{-3/2} \exp(-f/f_0), \quad f_0 = 0.3 \text{ Hz}$$

⇒ The total observed spectrum depends on the contribution (percentage) of each event.

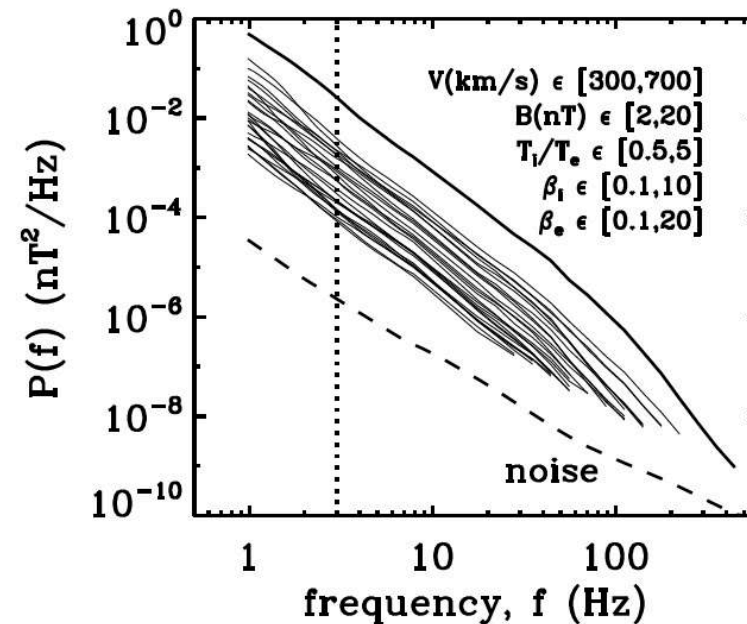
⇒ These results may explain spectral variability around ion scales.

[S. Lion, O. Alexandrova, A. Zaslavsky, 2015, submitted to APJ]

Turbulence at kinetic scales

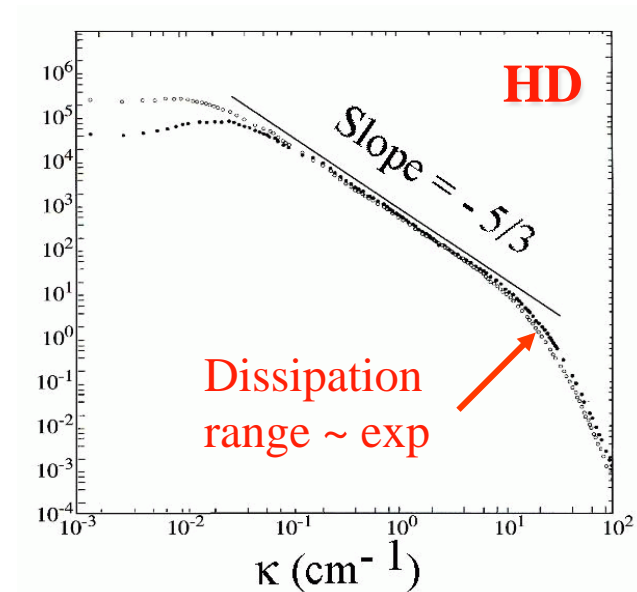
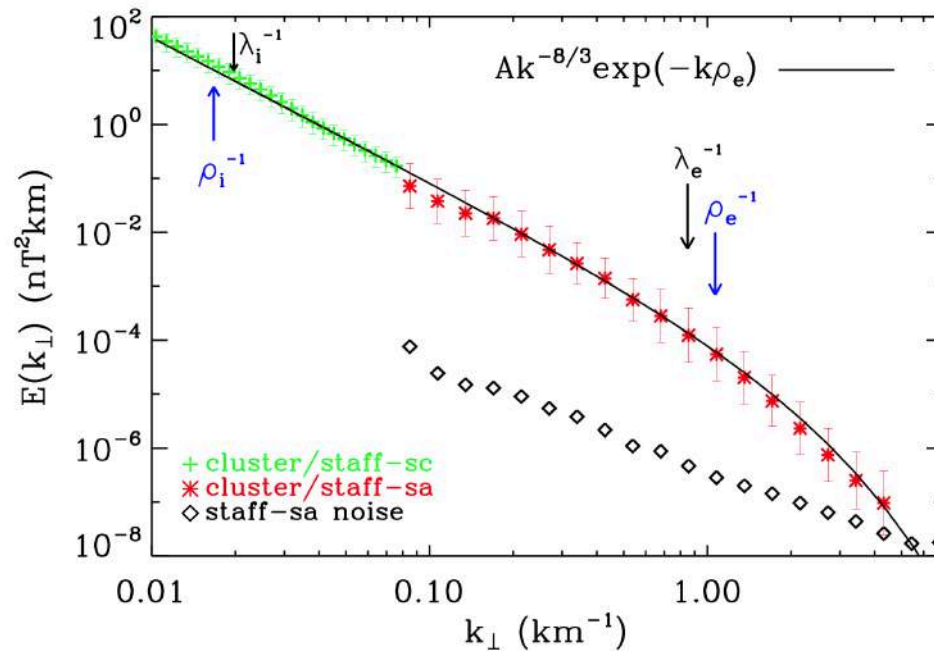
2. Electron scales

Cluster mission : the most sensitive instrumentation (magnetic spectrum up to 400 Hz, i.e. scales $\sim 300\text{m}$).



Turbulent spectrum at electron scales: dissipation range?!

[Alexandrova et al., 2012, APJ]



- In HD turbulence, dissipation range can be described by [Chen et al., 1993, PRL] :

$$E(k) = Ak^{-\alpha} \exp(-k/k_d)$$

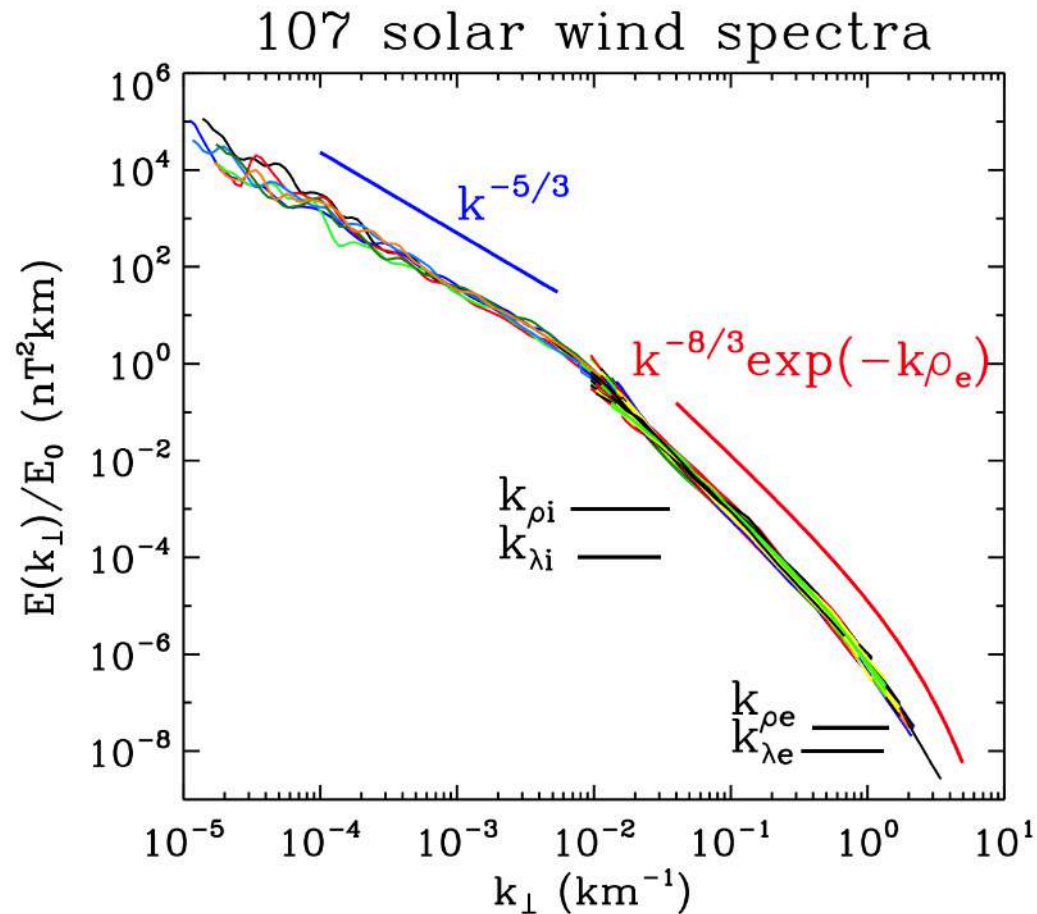
- In solar wind turbulence, we find a similar law :

$$E(k) = Ak^{-8/3} \exp(-k\rho_e)$$

General spectrum at electron scales

[Alexandrova et al., 2012, APJ]

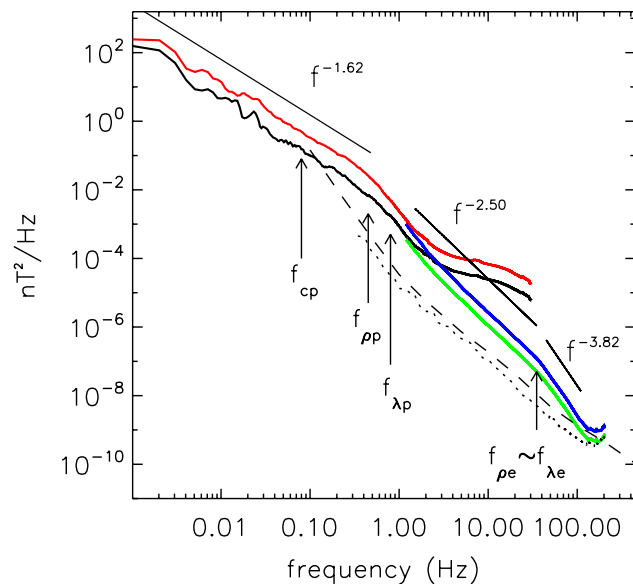
$$E(k) = Ak^{-8/3} \exp(-k\rho_e)$$



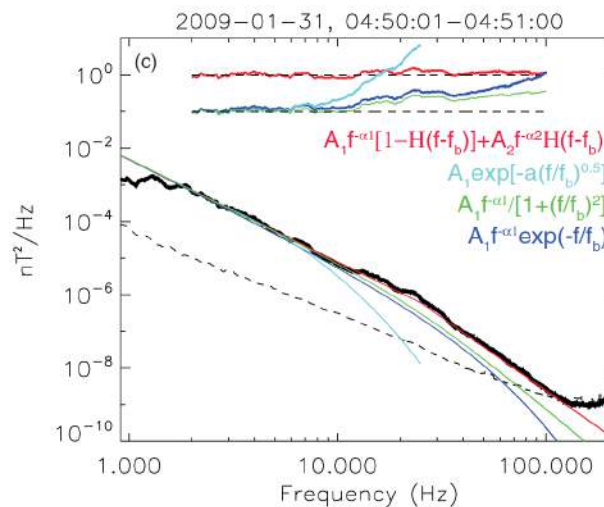
Electron Larmor radius ρ_e plays the role of the dissipation scale in collisionless solar wind turbulence.

Examples of different (non-universal) spectra at electron scales

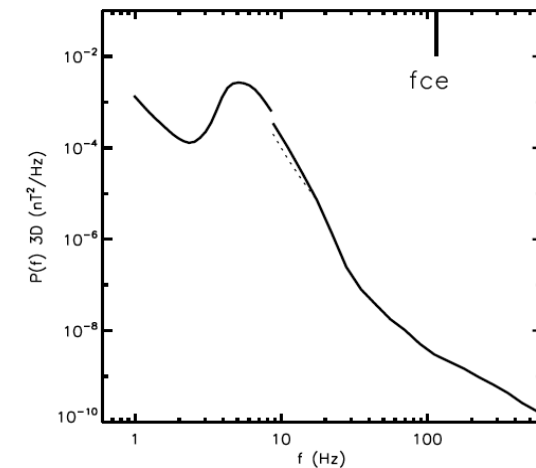
Observation of spectral break of bosse at electron scales:



[Sahraoui et al., 2009]



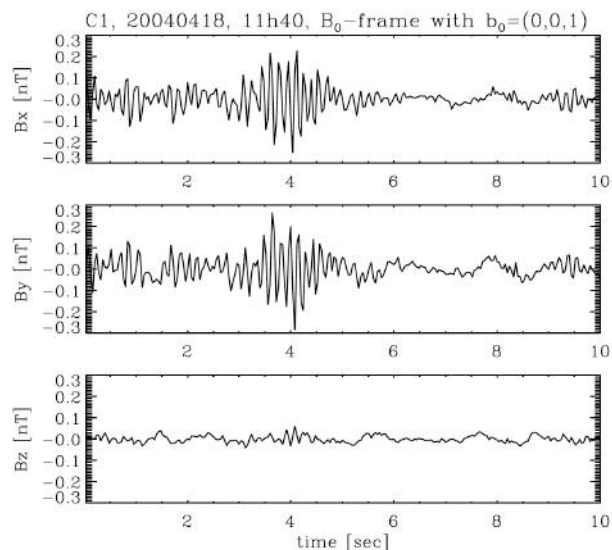
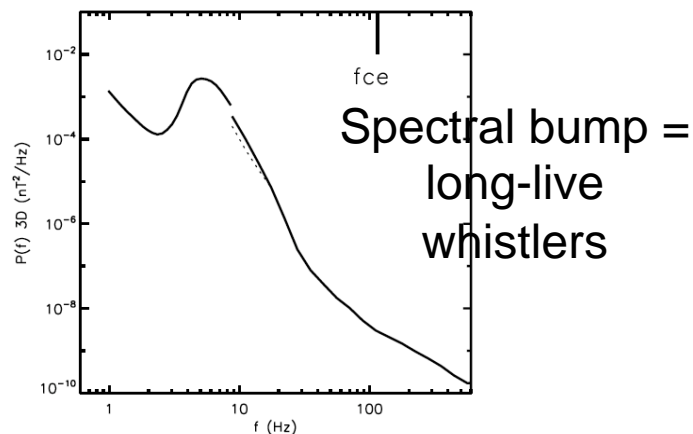
[Sahraoui et al., 2013]



[Lacombe et al. 2014]

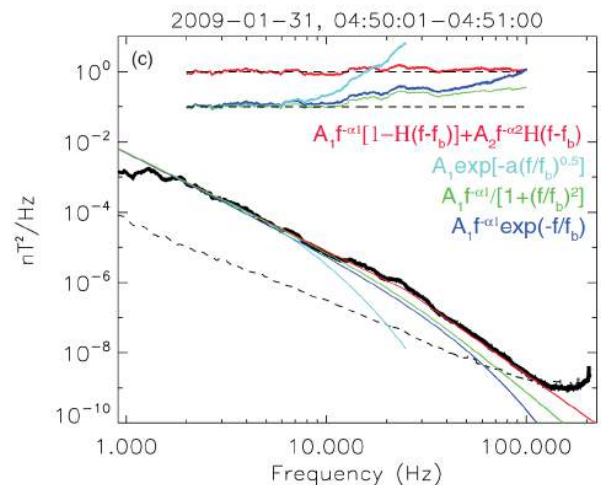
These non-universal features are due to appearance of quasi-monochromatic whistler waves in parallel propagation (with $k \parallel B_0$).

Examples of quasi-monochromatic whistler waves in the solar wind (5-10% of data)

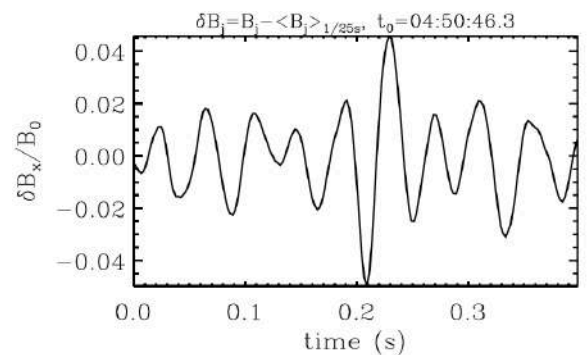


Phase-diff. between Bx and By:

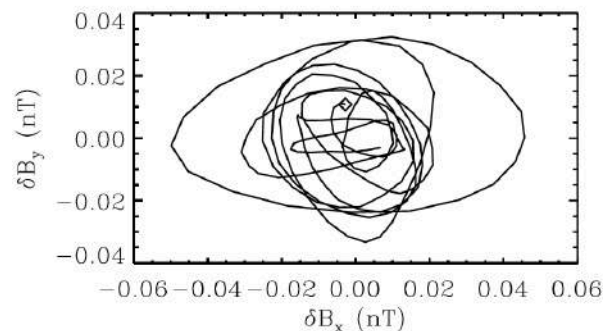
$$\Delta\Phi_{xy} = \Phi_x - \Phi_y = 90^\circ$$



[Sahraoui et al. 2013]: spectral break is due to sporadic whistlers



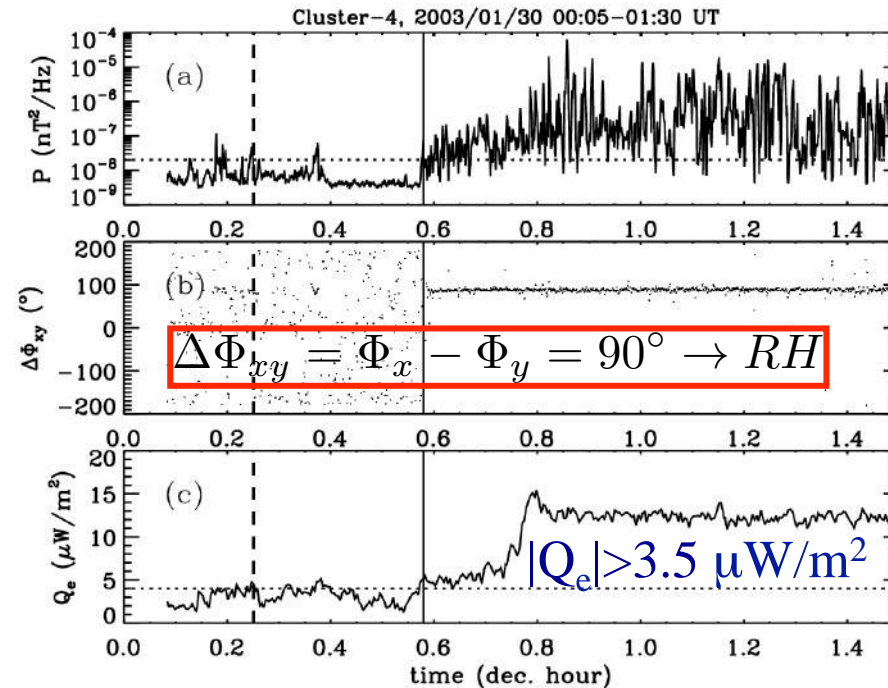
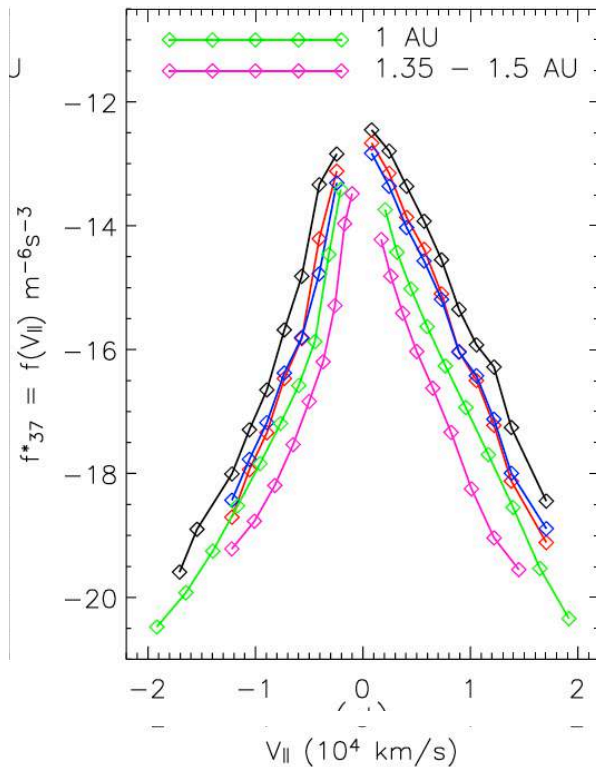
[Lacombe et al. 2014]



Whistlers appear for high electron heat flux

[Lacombe et al. 2014, APJ]

[Maksimovic et al. 2005]



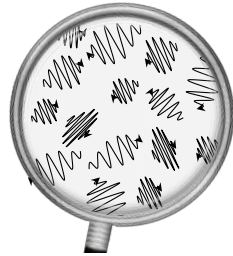
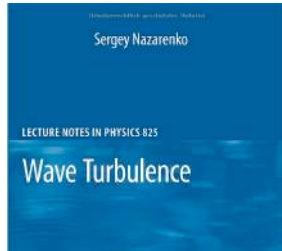
$$Q_e = \int \frac{m}{2} \mathbf{U} \mathbf{U}^2 f(v) d^3v, \quad \mathbf{U} = \mathbf{v} - \langle \mathbf{v} \rangle$$

Electron heat flux, Q_e , is a measure of the asymmetry of the electron distribution function $f(v_e)$. In the solar wind it is present for $f(v_{e||})$.

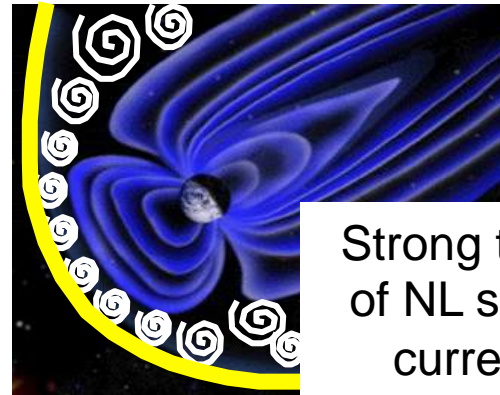
We find that whistlers grow with Q_e (generation mechanism: whistler heat flux instability, see Gary et al., 99).

Turbulence nature: weak (or wave) vs strong

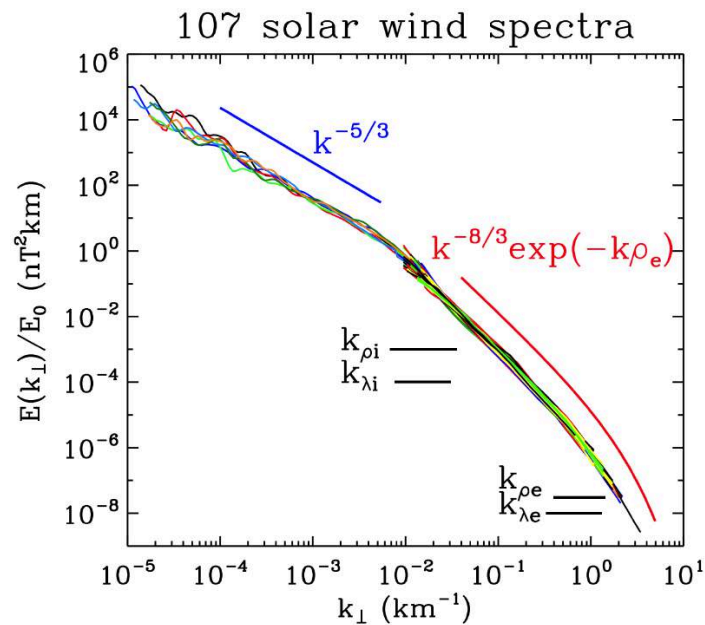
Courtesy S. Galtier



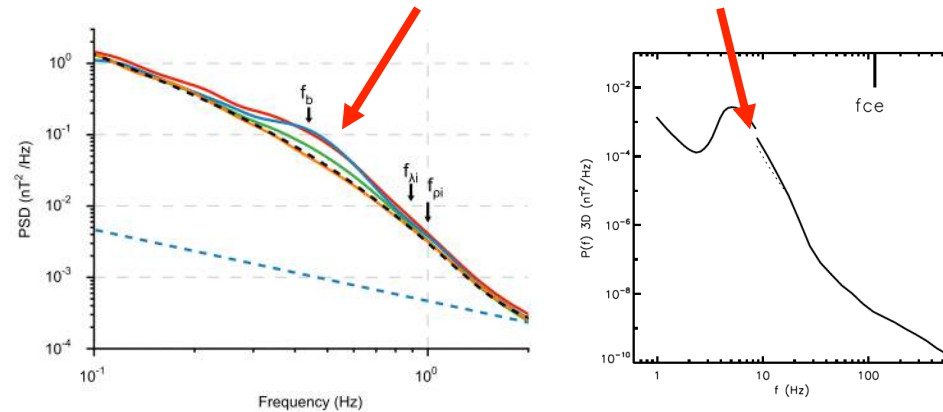
Weak turbulence:
mixture of waves with
+/-random phases



Strong turbulence: mixture
of NL structures (vortices,
current sheets, ect...)



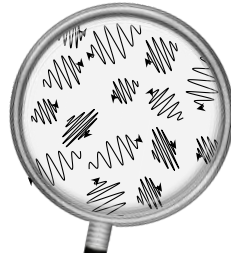
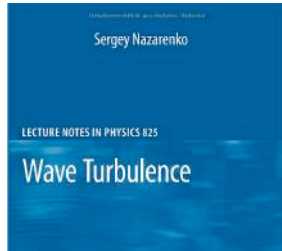
Zooms around ion and electron scales:



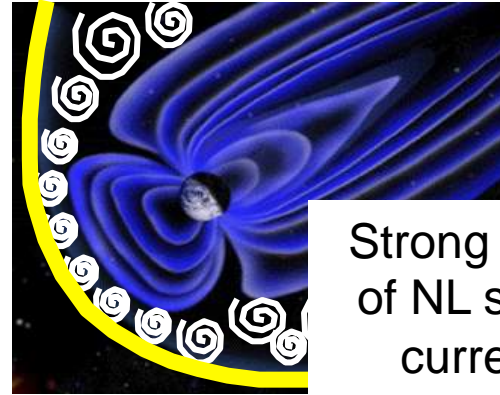
Strong turbulence within $k^{-5/3}$ and $k^{-8/3}$
ranges + waves at ion and electron
scales.

Turbulence nature: weak (or wave) vs strong

Courtesy S. Galtier

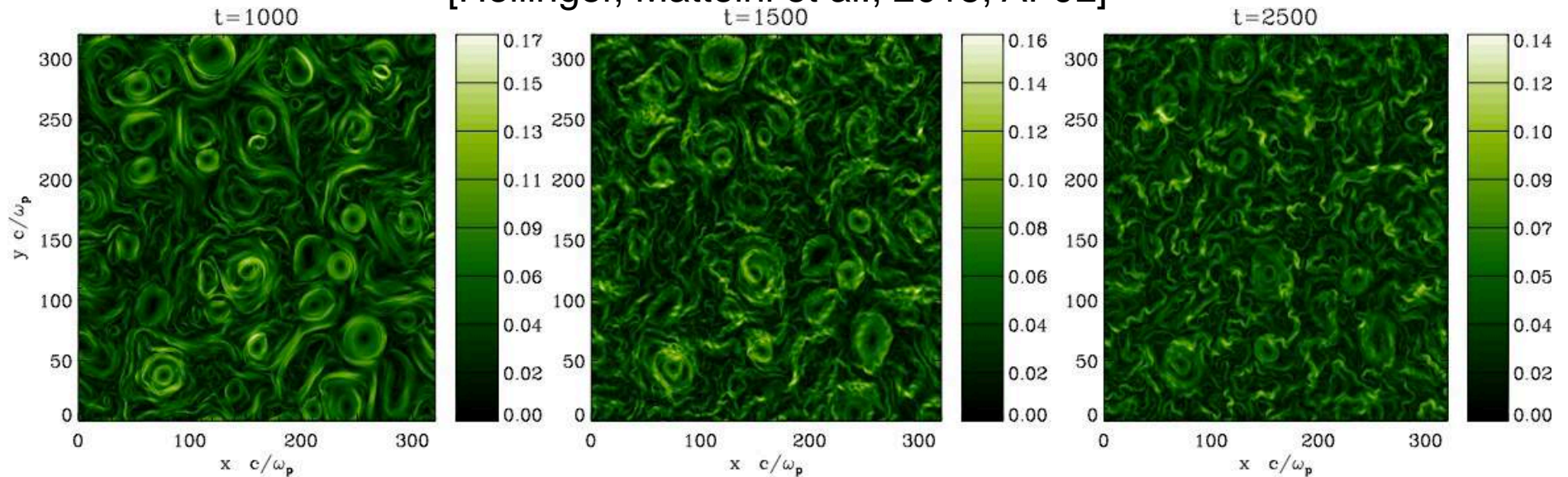


Weak turbulence:
mixture of waves with
+/-random phases



Strong turbulence: mixture
of NL structures (vortices,
current sheets, ect...)

[Hellinger, Matteini et al., 2015, APJL]

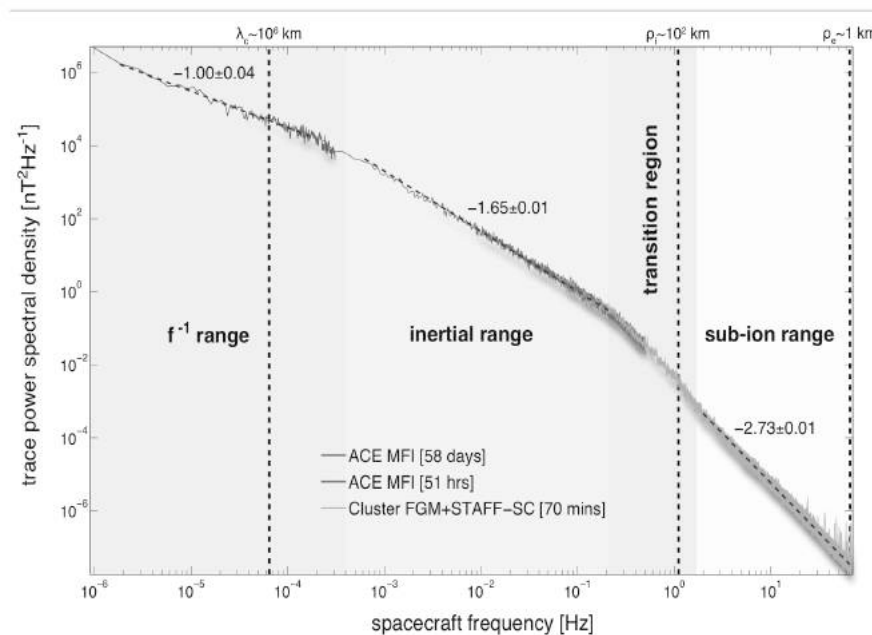


Courtesy of Lorenzo Matteini: 2D Hybrid numerical simulations showing development of strong turbulence (vortices) with superposed waves at ion scales.

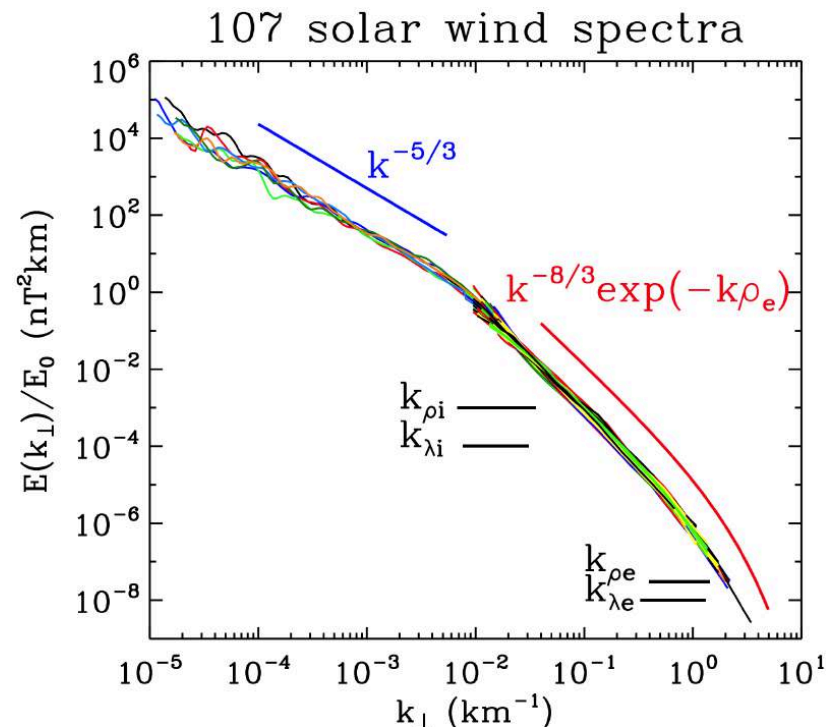
Discussion and Conclusion

- Solar wind is one of the best laboratories of space plasma turbulence.
- **In-situ observations: from MHD (10^7 km) to sub-electron scales (300 m).**
- At very large scales: spectrum is $\sim f^{-1}$, Alfvénic fluctuations;
- MHD inertial range is dominated by k_{\perp} fluctuations, with Kolmogorov's power law ;
- One decade around ion scales: superposition of turbulence and waves/ion instabilities ;
- Between ion and e-scales: small-scale inertial range $\sim k_{\perp}^{-8/3}$;
- One decade around e-scales: dissipation at $\sim \rho_e$ + sometimes whistler waves with k_{\parallel} .

Typical spectral density of solar wind turbulence



[Kiyani et al. 2015]

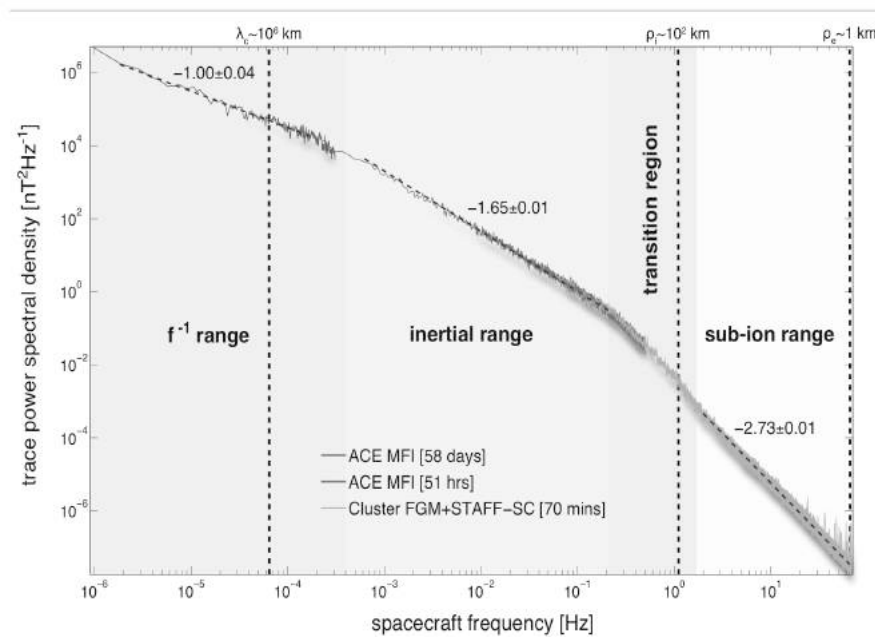


[Alexandrova et al. 2012]

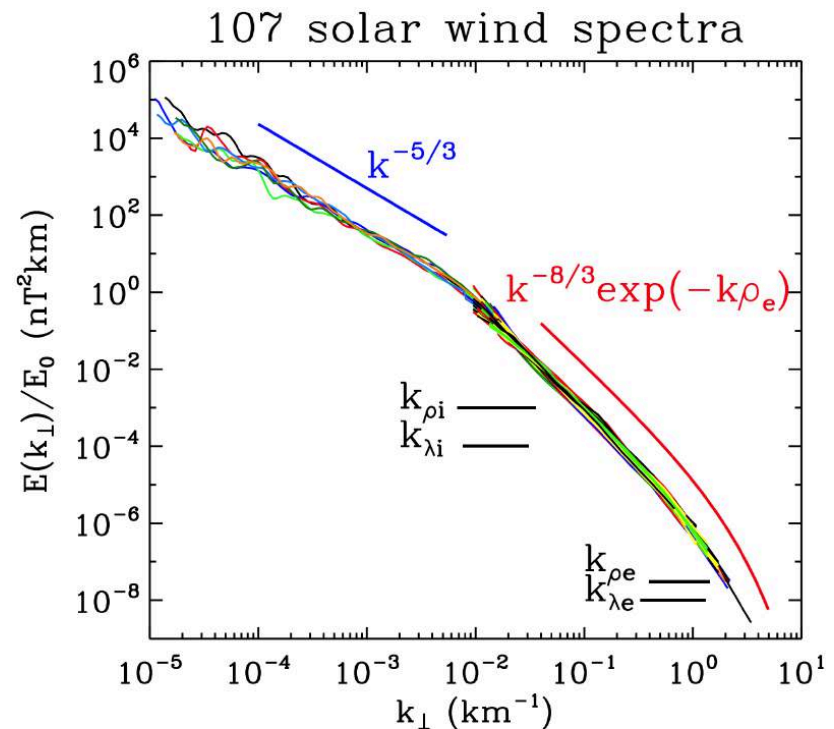
Open questions

- Nature of Kolmogorov like turbulence ? Role of compressibility ?
- Physical processes at ion scales ?
- Final dissipation at electron scales ?
- Plasma heating and particle acceleration by turbulence ?
- Dissipation without collisions ?
- ...

Typical spectral density of solar wind turbulence



[Kiyani et al. 2015]



[Alexandrova et al. 2012]

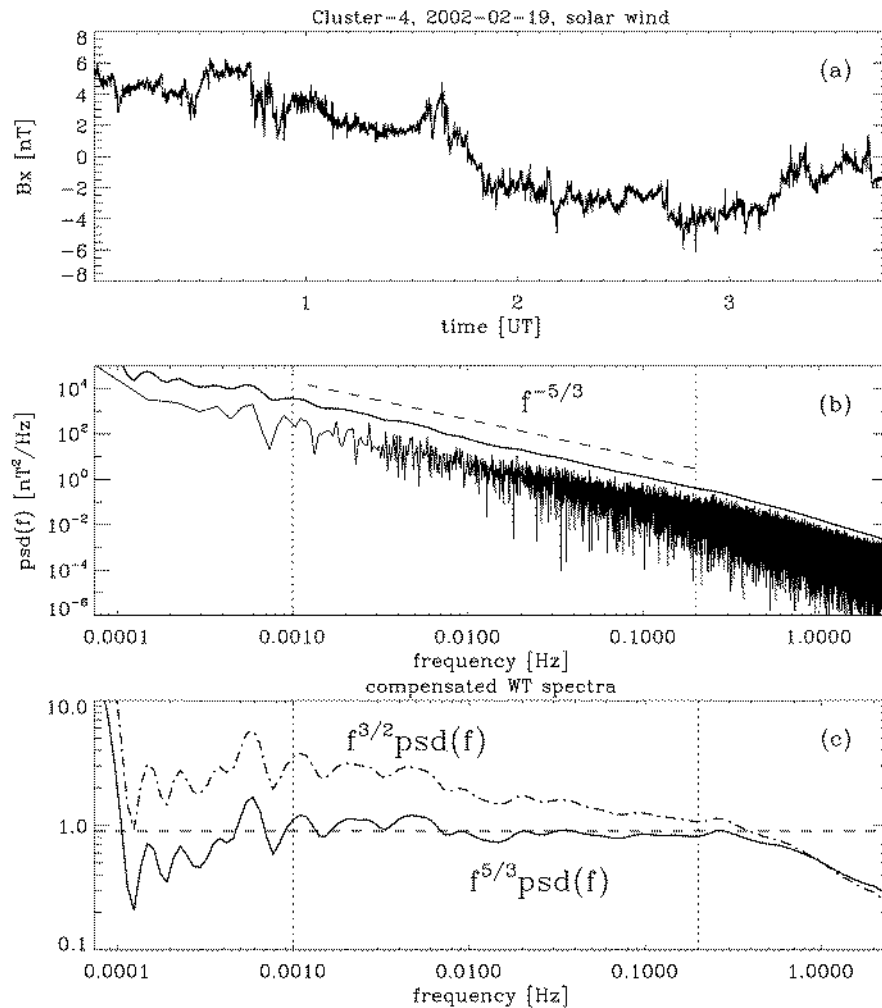
Bonus

How do we measure turbulent spectra?

Satellites in-situ measurements are time series =>

Fourier (or Wavelet) transform => frequency spectra

Methods for Characterising Microphysical Processes in Plasmas



- example of Cluster/FGM
(5 vectors/sec measurements)

How do we get k -
spectra?

Taylor hypothesis:

$$\ell = V_{sw}\tau = V_{sw}/f$$
$$k = 2\pi/\ell = 2\pi f/V_{sw}$$

[Dudok de Wit et al. 2013, SSR]

Taylor hypothesis

$$\omega_{obs} = \omega_0 + \mathbf{k} \cdot \mathbf{V}$$

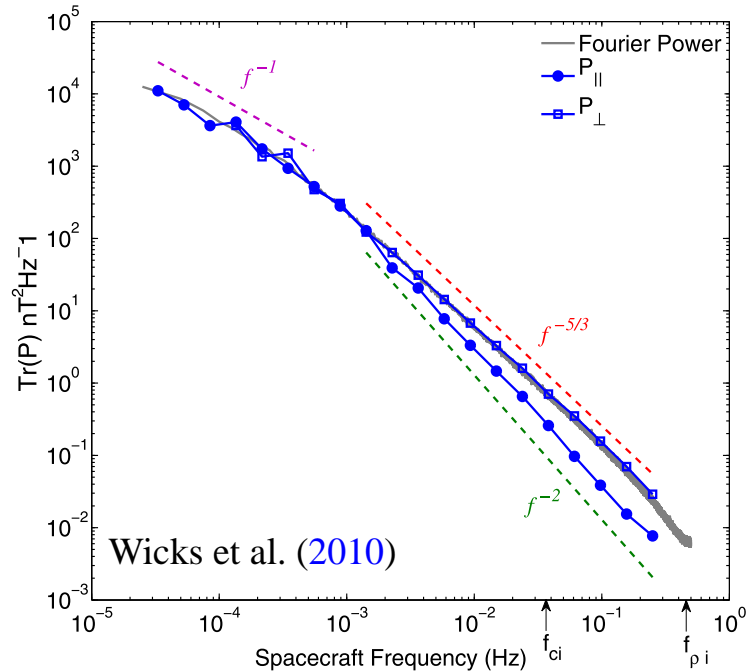
Supposing that $\omega_0 \ll k \cdot V$, ($V_\varphi \ll V$):

$$\omega_{obs} = \mathbf{k} \cdot \mathbf{V} = kV \cos(\Theta_{kV})$$

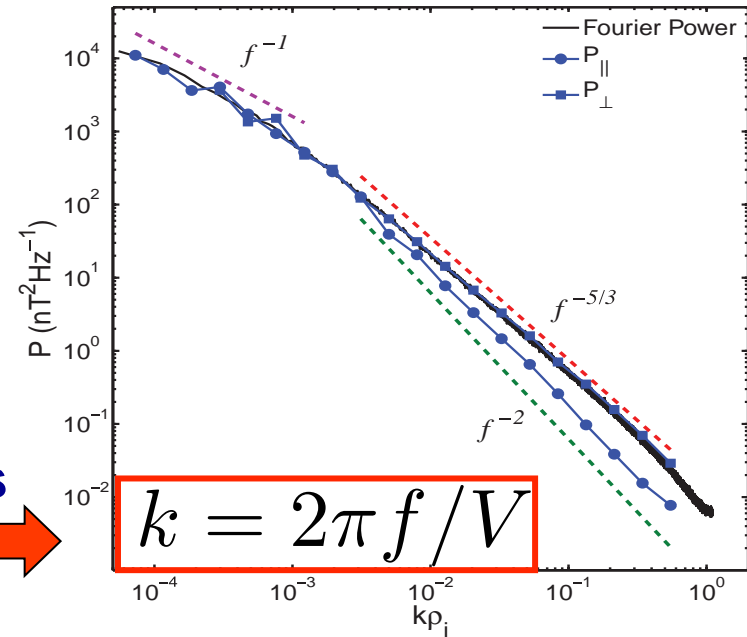
We don't know the angle between k and V
 \Rightarrow assumption of $k \parallel V$:

$$\omega_{obs} = kV \rightarrow k = 2\pi f / V$$

Anisotropy of turbulent fluctuations at MHD scales



Taylor hypothesis



- Alfvénic turbulence of Goldreich and Sridhar, 1995 is based on the idea of a balance between linear Alfvén time (along B_0) and non-linear time (in plane perp. to B_0):

$$\tau_A = \frac{\ell_{\parallel}}{V_A} \sim \tau_{NL} = \frac{\ell_{\perp}}{\delta V_{\perp}}$$

$$P(k_{\perp}) \sim k_{\perp}^{-5/3} ; P(k_{\parallel}) \sim k_{\parallel}^{-2}$$